

# LECTURE 3 -

## CHARGED INTERFACES

- RESOURCES:

- Phys. & Chem.

- OF I.F. [BUTT ET AL.]

- CHAPTER 4

3RD EDITION  
(2013)

- CHARGED MERCURY/WATER  
INTERFACE & THE  
LIPPMANN EQUATION

RELATIONS ENTRE LES PHÉNOMÈNES ÉLECTRIQUES  
ET CAPILLAIRES;

PAR M. GABRIEL LIPPMANN,

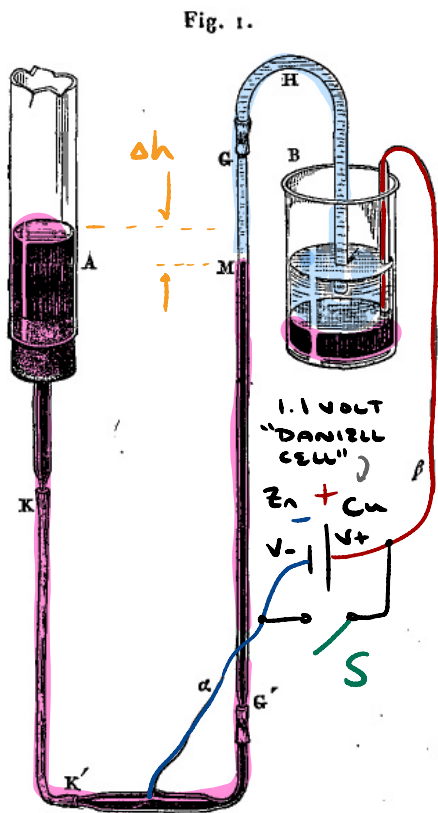
Ancien élève de l'École Normale supérieure.

ANN. CHIM. PHYS. 5<sup>e</sup> t. V  
1875, p. 494



- L'étude des forces électromotrices de contact a fourni jusqu'à présent un chapitre entièrement distinct de l'étude du phénomène capillaire ...
- On eût sans doute pensé à rapprocher l'une de l'autre ces deux propriétés physiques des surfaces de contact, la force électromotrice et la tension superficielle, et à chercher une relation constante entre elles, ...

- TO STUDY THE RELATIONSHIP BETWEEN ELECTRICITY & SURFACE TENSION, FRENCH SCIENTIST \* GABRIEL LIPPMANN (1845-1921) USED THE FOLLOWING APPARATUS:



- LIQUID MERCURY
- DILUTE ACID

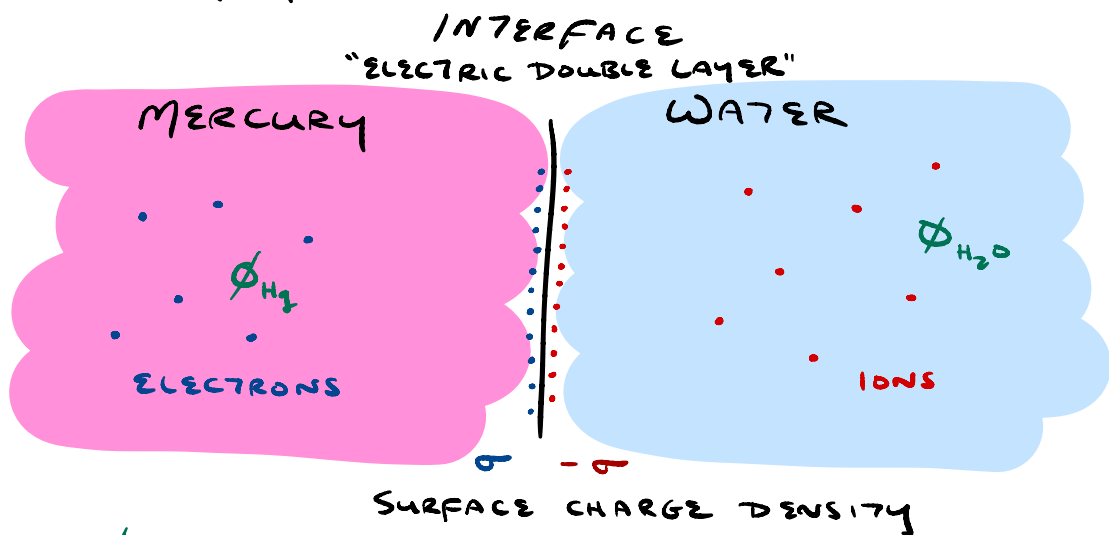
- W/ THE SWITCH  $S$  CLOSED, THE CAPILLARY FALL IS  $\Delta h = 14.0 \text{ mm}$

- QUIZ : IF THE RADIUS OF THE CAPILLARY IS  $R = 0.32 \text{ mm}$ , THEN WHAT IS THE IMPLIED SURFACE TENSION @ THE MERCURY/WATER INTERFACE? \*

- \* WHEN LIPPMANN OPENS THE SWITCH  $S$ , THE MENISCUS IS OBSERVED TO FALL ANOTHER  $4.9 \text{ mm}$  TO  $\Delta h = 18.9 \text{ mm}$ !

- A  $\approx 1V$  POTENTIAL APPLIED BETWEEN THE TWO VOLUMES OF MERCURY ARE ENOUGH TO CREATE A 35% INCREASE IN THE SURFACE TENSION!

- THERMODYNAMIC ANALYSIS :



$\phi$  : ELECTRIC POTENTIAL

- WE START W/ THE GIBBS ISO-THERM :

$$d\gamma = - \underbrace{\prod_{\text{ION}} d\mu_{\text{ION}}^s}_{\text{light blue bar}} - \underbrace{\prod_{\text{elec}} d\mu_{\text{elec}}^s}_{\text{pink bar}}$$

- IN EQUILIBRIUM, THE CHEMICAL POT.  $\mu^s$  OF THE ELECTRONS AND IONS @ THE INTERFACE ARE EQUAL TO THE CORRES-

PONDING CHEM. POTS  $\mu^{\text{BULK}}$  OF THE BULK PHASES.

- BECAUSE WE DEAL NOW W/ CHARGED PARTICLES, WE MUST INCLUDE THE ELECTRIC POTENTIAL  $\phi$  IN THE CHEMICAL

POTENTIAL:

$$\left[ F = e N_A \approx 9.6 \text{ C/mol} \right]$$

$$\mu_{\text{elec}}^{\text{BULK}} = \mu_{\text{elec}}^{\circ} - F \phi_{\text{Hg}} \quad \text{"FARADAY CONSTANT"}$$

$$\mu_{\text{ION}}^{\text{BULK}} = \mu_{\text{ION}}^{\circ} + F z \phi_{\text{H}_2\text{O}}$$

ION CHARGE

- IF WE CONSIDER ONLY CHANGES DUE TO THE APPLIED VOLTAGE  $V_+ - V_-$ , THEN WE CAN IGNORE  $\mu_{\text{elec}}^{\circ}$  &  $\mu_{\text{ION}}^{\circ}$ :

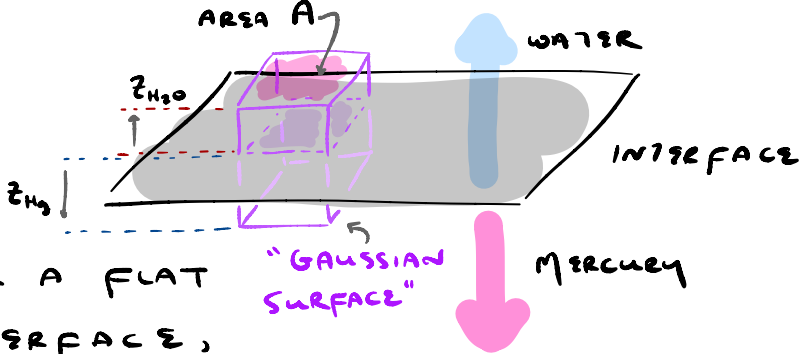
$$d\mu_{\text{elec}}^{\text{S}} = d\mu_{\text{elec}}^{\text{BULK}} = -F d\phi_{\text{Hg}}$$

$$d\mu_{\text{ION}}^{\text{S}} = +F z d\phi_{\text{H}_2\text{O}}$$

• "ELECTRON NEUTRALITY":

- BOTH MERCURY AND THE WEAK ACID SOLUTION ARE ELECTRICAL CONDUCTORS.

- THEREFORE IN THE BULK PHASES WE REQUIRE THE ELECTRIC FIELD  $\vec{E} = 0$ .



- FOR A FLAT INTERFACE,

WE HAVE FROM GAUSS' LAW:

$$\oint \vec{D} \cdot d\vec{A} = \int dV \rho = Q_{ENC}^*$$

$$\left[ \vec{D}(z) = \epsilon_r(z) \epsilon_0 \vec{E}(z) \right]$$

"DISPLACEMENT FIELD"

$$\epsilon_r^{H_2O} \epsilon_0 A E_z(z_{H_2O}) + \epsilon_r^{H_2O} \epsilon_0 A E_z(z_{H_2O}) = A \sigma_{elec} + A \sigma_{ion}$$

$$\Rightarrow \sigma_{elec} = - \sigma_{ion}$$

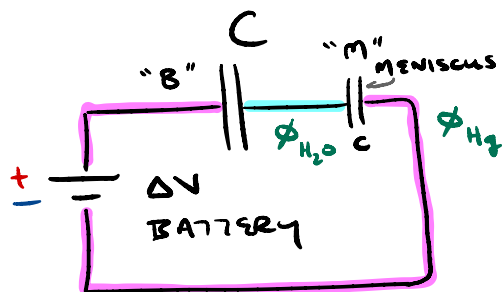
"ELECTRO-NEUTRALITY"



- BUT:  $\sigma_{\text{elec}} = -F \Gamma_{\text{elec}}$   
 $\sigma_{\text{ion}} = +zF \Gamma_{\text{ion}}$

$$\begin{aligned} \Rightarrow d\gamma &= -\Gamma_{\text{elec}} \times \left( -F d\phi_{\text{Hg}} \right) \\ &\quad - \Gamma_{\text{ion}} \left( +Fz d\phi_{\text{H}_2\text{O}} \right) \\ &= +\sigma d \left( \phi_{\text{H}_2\text{O}} - \phi_{\text{Hg}} \right) \end{aligned}$$

- NOW WE CAN MODEL LIPPMAN'S EXPERIMENT AS THE FOLLOWING CIRCUIT:



- BOTH MERCURY/WATER INTERFACES "B" + "M" WILL ACT LIKE CAPACITORS w/ IDENTICAL CAPACITANCES PER UNIT AREA.

- BECAUSE THE MENISCUS IS MUCH SMALLER, IT'S CAPACITANCE  $C$  WILL BE MUCH SMALLER THAN THE INTERFACE @ "B" CAPACITANCE  $C$ .

• THE CHARGE ON THE MENISCUS IS  
 $\pm \sigma a$  <sup>✓ SURF. AREA OF MENISC.</sup>, AND THIS MUST ALSO BE THE  
 CHARGE ON THE LARGER INTERFACE.

• THE VOLTAGE ACROSS THE MENISCUS IS

$$\text{THUS } V_M = \frac{\sigma a}{C}$$

• AND ACROSS THE OTHER:

$$V_B = \frac{\sigma a}{C}$$

• SO  $\frac{V_M}{V_B} = \frac{C}{C} \gg 1$

• SO:

$\approx$  BATTERY  
VOLTAGE

$$d(\phi_{H_2O} - \phi_{Hg}) = dV_M = d\Delta V$$

$$\Rightarrow d\phi = \sigma d\Delta V$$

$$\Rightarrow \left[ \frac{d\phi}{d\Delta V} = \sigma \right] \left[ \text{WHERE } \Delta V = V_{H_2O} - V_{Hg} \right]$$

"LIPPMANN EQN"

VOLTAGE  
Applied  
To  $H_2O$  -  
SIDE OF  
MENISCUS

VOLTAGE  
Applied  
To  $Hg$  -  
SIDE OF  
MENISCUS

• BUT WE ALSO HAVE:

\*  $Q = -a\sigma$

$$\Rightarrow \left[ C_A \equiv -\frac{1}{a} \frac{dQ}{d\Delta V} = -\frac{\cancel{a}}{\cancel{a}} \frac{d\sigma}{d\Delta V} = -\frac{d^2\phi}{d\Delta V^2} \right]$$

\* Q HERE REFERS  
 TO CHARGE ON  
AQUEOUS SIDE  
 OF THE INTERFACE.

"DIFFERENTIAL CAPACITANCE PER UNIT AREA"

# • SURFACE CHARGE AND $\gamma$ vs. $\Delta V$ CURVES.

- ANOTHER FRENCH PHYSICIST J. G. Gouy [1854 - 1926] PERFORMED CAREFUL MEASUREMENTS OF  $\gamma$  vs.  $\Delta V$  FOR THE MERCURY/AQUEOUS INTERFACE FOR A LARGE NUMBER OF SOLUTIONS & CONCENTRATIONS:

ANN. CHEM. PHYS. [7] 29, 145 (1903)

M. Lippmann, à la suite d'expériences peu nombreuses, avait cru que la fonction électrocapillaire était la même, quelle que fût la composition chimique du liquide, et que la dérivée  $\frac{d^2\theta}{d\Delta^2}$  était une constante <sup>(1)</sup>; dès lors toutes les courbes électrocapillaires étaient une même parabole. Dès le début de mes recherches, j'eus l'occasion de constater que les courbes électrocapillaires sont bien différentes les unes des autres, soit par la valeur du maximum, soit par leur forme et leur dissymétrie plus ou moins marquée <sup>(2)</sup>. Chaque liquide a donc sa courbe, électrocapillaire propre, et leur étude forme un sujet des plus étendus, auquel ce travail est consacré en grande partie.

Gouy

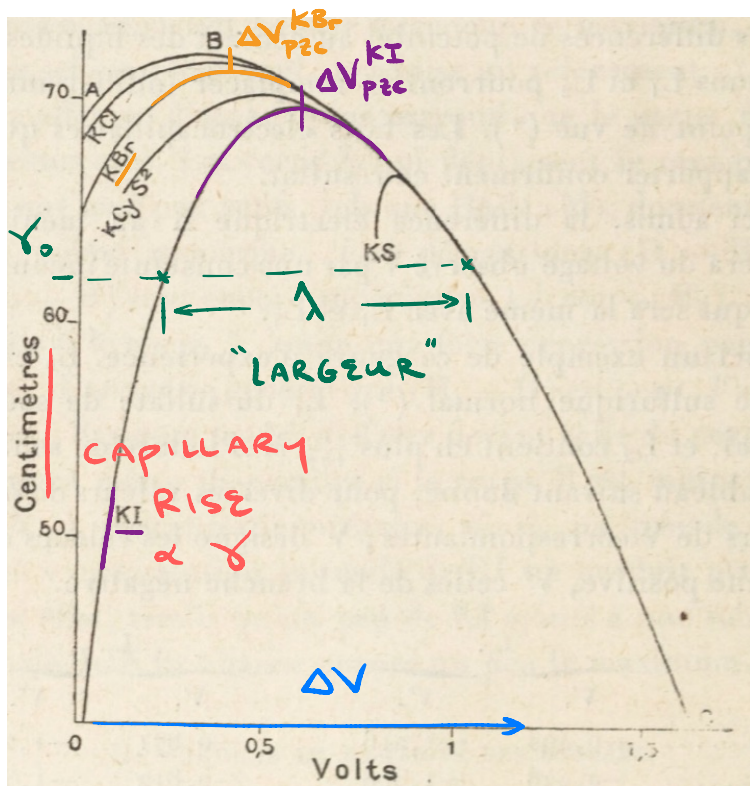


$$[\theta = \gamma, \Delta = \Delta V]$$

- FOR EXAMPLE, H<sub>2</sub> MEASURED  $\gamma$  vs.  $\Delta V$  FOR SEVERAL POTASSIUM SALTS:

III. — APPAREILS.

La méthode de l'électromètre capillaire, que j'ai employée, ne diffère pas en principe de celle dont M. Lippmann a fait usage. Diverses modifications ont été apportées aux appareils.



- WHAT WE MIGHT FIRST NOTE:
  - THE PEAKS OF THE CURVES DIFFER FOR DIFFERENT ANIONS

• E.G.  $\frac{d\gamma}{d\Delta V} (\Delta V = 0.5V) \neq 0$

WHILE  $\frac{d\gamma}{d\Delta V} (\Delta V = 0.3V) \neq 0$

BUT:  $\frac{d\gamma}{d\Delta V} = \sigma^*$

\* RECALL: SURFACE CHARGE IS DEFINED TO BE THAT OF THE METAL SIDE OF THE INTERFACE.

SO: IONS SPONTANEOUSLY ADSORB ONTO MERCURY/AQUEOUS INTERFACE.

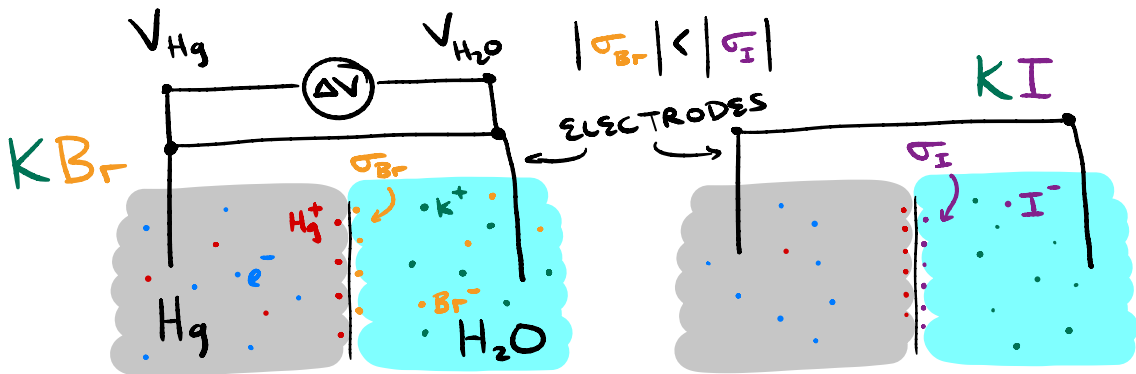
"SPECIFIC ADSORPTION"

• I.E. THE "POINT OF ZERO CHARGE" [PZC]  $\Delta V_{PZC}$

WHERE  $\frac{d\gamma}{d\Delta V} (\Delta V_{PZC}) = 0$

$$\Delta V = V_{H_2O} - V_{Hg} = 0$$

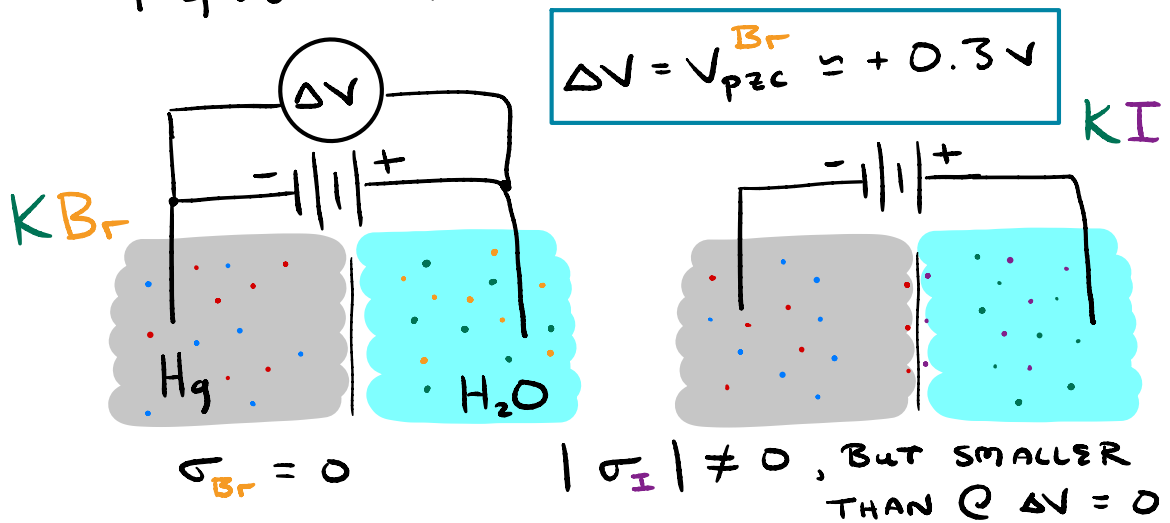
DEPENDS ON THE SOLUTION:

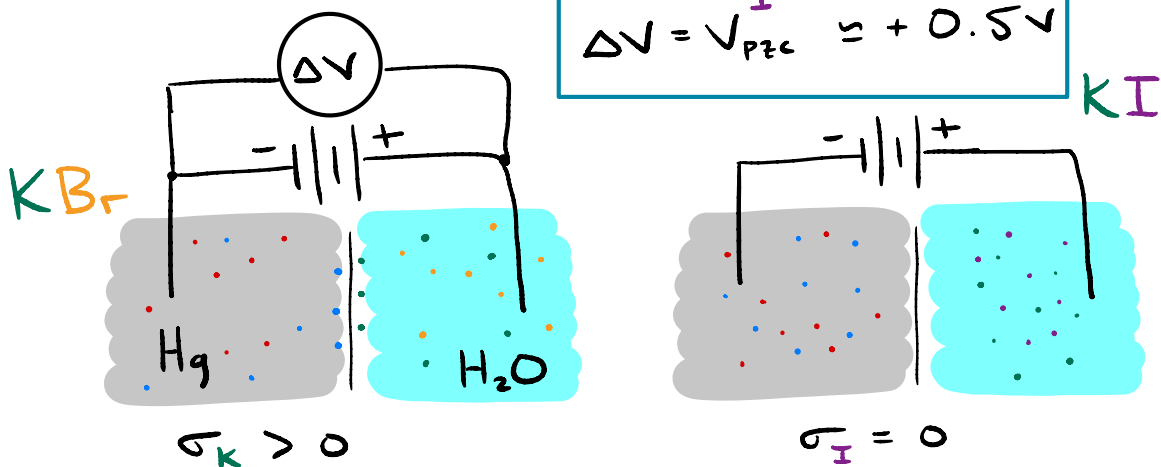


• @  $\Delta V = 0$  :  $-\frac{d\gamma^{KI}}{d\Delta V} = \sigma_I < \sigma_{Br} < 0$

• TO DEPLETE THIS NEGATIVE SURFACE CHARGE, WE APPLY A POSITIVE VOLTAGE  $\Delta V = V_{H_2O} - V_{Hg}$  BETWEEN THE AQUEOUS SOLUTION & THE MERCURY.

• THE POSITIVE CHARGE ON THE AQUEOUS ELECTRODE ATTRACTS THE ANIONS AWAY FROM THE INTERFACE, LIKEWISE THE NEGATIVE CHARGE ON THE MERCURY ELECTRODE REPELS THE ANIONS :





- WHY DO ANIONS BIND TO MERCURY [AND OTHER METALS (Au, Pt, Ag)] MORE THAN CATIONS?
- ACCORDING TO P.C.I.: IT'S COMPLICATED. BUT PART OF THE ANSWER LIES IN THE STRONGER "HYDRATION" OF CATIONS:
  - THE COST OF "SHEDDING" THE ORDERED SHELL OF  $H_2O$  MOLECULES IN ORDER TO BIND TO METAL SURFACE IS TOO HIGH.

• "

## II. — LARGEUR DES COURBES.

"

9

DOUBLE LAYER  
CAPACITANCE.

- Gouy ALSO STUDIED THE VOLTAGE DIFFERENCE ["LARGEUR"]

$\lambda$  BETWEEN TWO POINTS ON A  $\gamma$  vs.  $\Delta V$  CURVE OF EQUAL

REFERENCE SURFACE TENSION

$\gamma_0$  [SEE EXAMPLE CURVE ABOVE].

- BELOW WE SHOW THE RESULTS FOR  $\lambda$  FOR DIFFERENT  $\gamma_0$ .

FOR SOLUTIONS OF POTASSIUM  
NITRATE FOR DIFFERENT  
CONCENTRATIONS:

On y a inscrit, pour les hauteurs 940, 900, 860 et 820, les largeurs  $\lambda$  des courbes, et les accroissements  $\delta\lambda$  que prennent ces largeurs quand la concentration devient 10 fois plus petite.

CONCENTRATION

$1M \equiv 1g/liter$

TABLEAU V.

		940.		900.		860.		820.
	$h.$	$\lambda.$	$\delta\lambda.$	$\lambda.$	$\delta\lambda.$	$\lambda.$	$\delta\lambda.$	$\lambda.$
K Az O <sub>3</sub>	( $\frac{1}{100}$ M)	0,933		1,211		"		"
id.	( $\frac{1}{10}$ M)	0,862	0,071	1,149	0,062	1,377	"	"
id.	(1 M)	0,767	0,095	1,059	0,090	1,296	0,081	1,497

\* SURFACE TENSION MEASURED AS CAPILLARY FALL OF MENISCUS, NORMALIZED TO PURE WATER SOLUTION, SET EQUAL TO 1000.

• SO GOUY MEASURES A SYSTEMATIC DECREASE IN  $\gamma$  w/ INCREASING "KAZO<sub>3</sub>" CONCENTRATION.



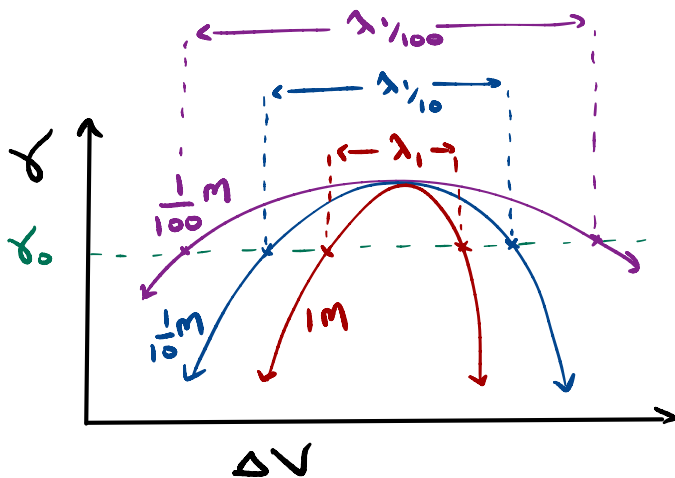
• WHAT DOES THIS IMPLY?

• NEAR MAXIMUM OF  $\gamma$  vs.  $\Delta V$  CURVE, WE HAVE:

\* APPROX. CONST. w/ CONCENTRATION

$$\gamma(\Delta V_{pzc} + \delta V) \approx \gamma(\Delta V_{pzc}) + \frac{1}{2} \frac{d^2\gamma}{d\Delta V^2}(\Delta V_{pzc}) \delta V^2$$

[ TAYLOR EXPANSION OF  $\gamma(\Delta V)$  ABOUT  $\Delta V_{pzc}$ . ]  
NOTE THAT  $d\gamma/d\Delta V(\Delta V_{pzc}) = 0$



- A smaller  $\lambda$  we see then corresponds to a larger CURVATURE OF THE PARABOLA, I.E. A LARGER  $\left| \frac{d^2 \phi}{d \Delta V^2} \right|$ .
- BUT WE KNOW  $-\frac{d^2 \phi}{d \Delta V^2} = C_A$
- SO INCREASING SOLUTE CONCENTRATION INCREASES THE CAPACITANCE OF THE DOUBLE LAYER!
- WHY?
- GOUY ALONG W/ ENGLISH CHEMIST D.L. CHAPMAN [1869-1958] WORKED OUT THE THEORY OF THE DIFFUSE LAYER TO EXPLAIN THESE [ & OTHER ] OBSERVATIONS.